Exploring turbulence through particles: statistical properties of spheroids in turbulent flows

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16/05/2023 Bangalore, ICTS International Centre for Theoretical Sciences

My institute

LILLE NORTHERN FRANCE: CONNECTED TO WHOLE EUROPE



74000 students 1800 PhD



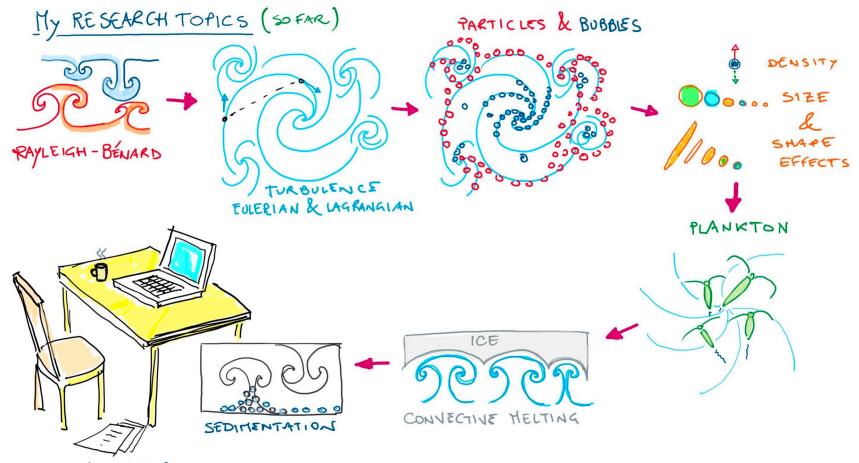




uml.univ-lille.fr

20 faculty, 12 PhD, 2 teams: a) complex fluids (numerics, theory) b) solids and structures (numerics, experiments)

My research in a doodle



www.ecalzavarini.info

Talk's outline

Introduction

- Recap on *models, simulations & experiments* on particles in turbulence
- Interesting questions

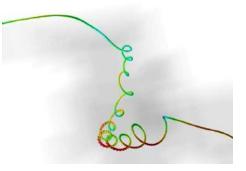
New results

- Our numerical study on *large neutral (a)spherical* particles in turbulence
- Acceleration statistics
- Rotation statistics
- Conclusions

A brief recap: Models for the motion of a particle in a flow (1)

• Fluid tracers, or the Lagrangian point of view

$$\frac{d\mathbf{x_f}}{dt} = \mathbf{u}(\mathbf{x_f}(t), t) \qquad \mathbf{a}_f \equiv \frac{D\mathbf{u}}{Dt} = -\frac{1}{\rho_f}\nabla p + \nu\Delta \mathbf{u} + \mathbf{f}$$



• Small inertial particles, an equation with many authors and many names

A brief recap: Models for the motion of a particle in a flow (2)

• Corrections to the point-particle equation

Finite-size effect (Faxén laws 1922)

$$\begin{split} \mathbf{u} &\to [\mathbf{u}]_S \simeq \mathbf{u} + \frac{r_p^2}{6} \Delta \mathbf{u} & \mathbf{u} \to [\mathbf{u}]_V \simeq \mathbf{u} + \frac{r_p^2}{10} \Delta \mathbf{u} \\ & \text{in drag , history} & \text{in fluid acceleration , added mass} \end{split}$$

Finite-Reynolds effect (Shiller-Neumann 1935) wake drag empirical correlation

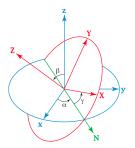
 $6\pi\mu r_p \to 6\pi\mu r_p (1+0.15Re_p^{0.687})$ $Re_p \lesssim 10^3$

Lift force for light particles (Auton JFM 1987, inviscid calculation)

$$\mathbf{f}_{lift} = \frac{m_f}{2} (\mathbf{u} - \mathbf{v}) \times (\nabla \times \mathbf{u})$$

• If particles are even larger: rigid body equations

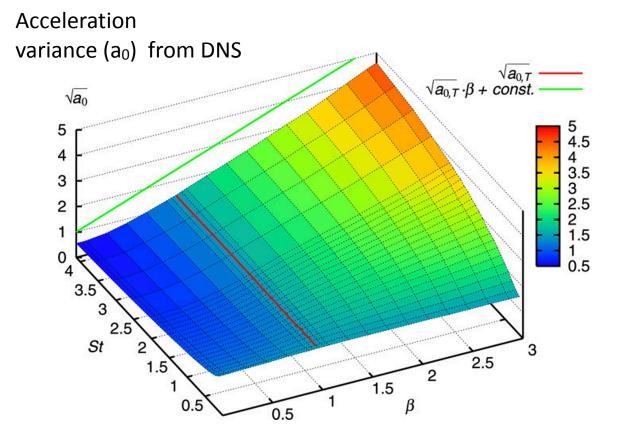
$$m_{p} \frac{\mathrm{d}\mathbf{v}}{\mathrm{d}t} = \mathbf{F} \qquad \mathbf{F} = \oint_{S_{p}} \boldsymbol{\sigma} \cdot \mathbf{n} \, dS$$
$$\frac{\mathrm{d}\boldsymbol{\mathcal{I}}\boldsymbol{\Omega}}{\mathrm{d}t} = \mathbf{T} \qquad \mathbf{T} = \oint_{S_{p}} (\mathbf{x} - \mathbf{r}) \times (\boldsymbol{\sigma} \cdot \mathbf{n}) \, dS$$



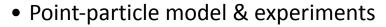
A brief recap: Acceleration statistics in numerics and experiments (1)

• Point-particle model & experiments

$$\frac{d\mathbf{v}}{dt} = \beta \frac{D\mathbf{u}}{Dt} + \frac{1}{\tau_p} (\mathbf{u} - \mathbf{v}) \qquad \frac{\beta = 3\rho_f / (\rho_f + 2\rho_p)}{\tau_p = r_p^2 / (3\nu\beta)}$$



A brief recap: Acceleration statistics in numerics and experiments (1)



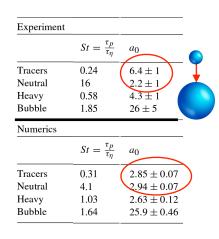
tracer d = 1.8 η



$$=\beta \frac{D\mathbf{u}}{Dt} + \frac{1}{\tau_n} (\mathbf{u} - \mathbf{v}) \qquad \beta = \frac{1}{\tau_n}$$

 $eta=3
ho_f/(
ho_f+2
ho_p)$ $au_p=r_p^2/(3
ueta)$

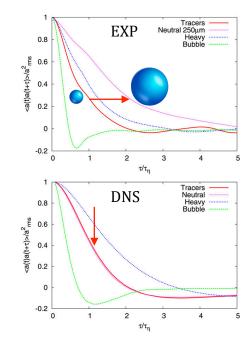
Variance (a₀)



Temporal correlation

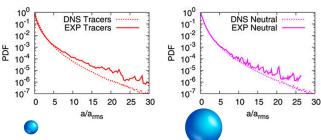
 $d\mathbf{v}$

 \overline{dt}



Probability density function

Volk, EC, Verhille, Lohse, Mordant, Pinton, Toschi, Physica D (2008)



A brief recap: Acceleration statistics in numerics and experiments (2)

EC, Volk, Bourgoin, Leveque, Pinton, Toschi, JFM (2009)

EXP D/η=12 EXP D/η=16 EXP D/η=25

DNS tracer

DNS D/ŋ=8

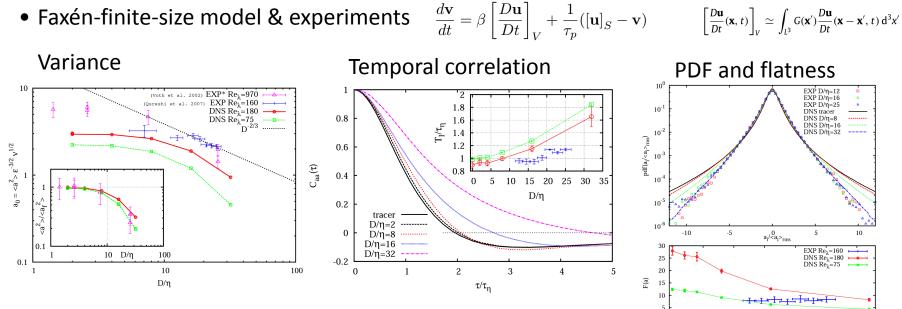
DNS D/n=16

DNS D/n=32

5

¹⁵ D/n 20 DNS Re₂=75

25



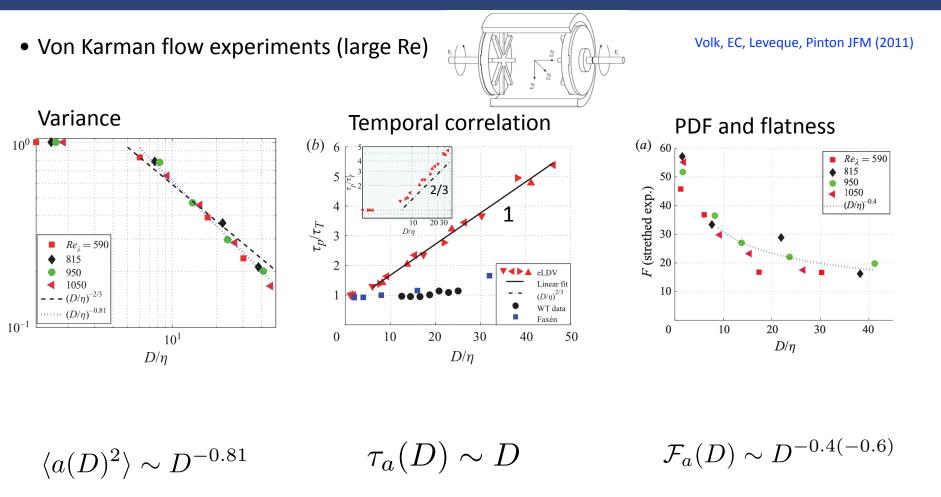
The coarse grained turbulence hypothesis (CGH)

• Faxén-finite-size model & experiments

 $\tau_u(\ell) \sim \frac{\ell}{\delta_\ell u} \sim \ell^{2/3}$ K41 $\delta_\ell u \sim \ell^{1/3}$ $\langle (\delta_\ell u)^p \rangle \sim \ell^{\zeta_p}$

9

A brief recap: Acceleration statistics in numerics and experiments (3)



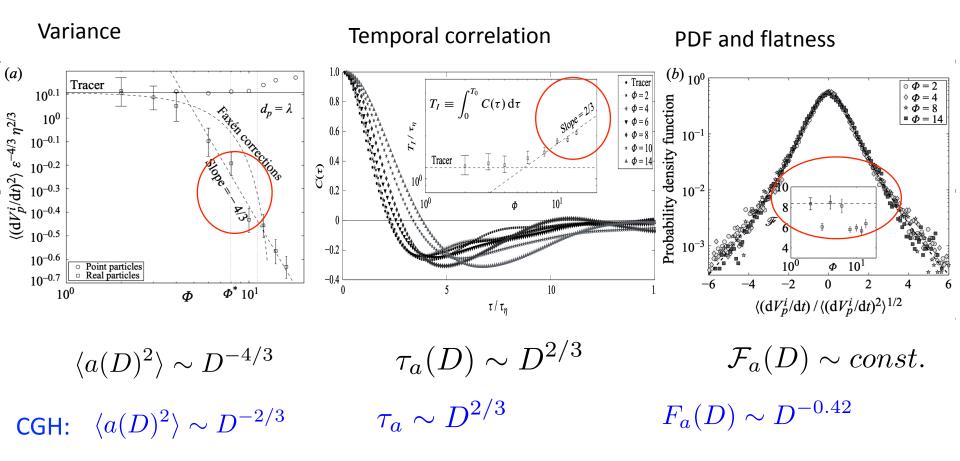
CGH: $\langle a(D)^2 \rangle \sim D^{-2/3}$

 $au_a \sim D^{2/3}$

 $F_a(D) \sim D^{-0.42}$

A brief recap: Acceleration statistics in numerics and experiments (4)

• Fully resolved simulations (low Re) $Re_{\lambda} = 32$



Cisse, Homann & Bec JFM (2013) -> slipping motion and particle boundary layer Uhlmann & Chouippe JFM (2017) -> particle preferential concentration Homann & Bec JFM (2010)

Interesting questions

• Do finite-sized particles experience coarse-grained turbulent flow accelerations?

$$a(D) \sim \frac{\delta_D p}{D} \sim \frac{(\delta_D u)^2}{D} \sim D^{-1/3} \qquad \tau_a(D) \sim \frac{D}{\delta_D u} \sim D^{2/3} \qquad \qquad \mathcal{F}_a(D) \sim D^{\zeta_8 - 2\zeta_4}$$

or does particle-flow coupling create specific statistical signatures?

• How is the particle finite-size reflected on particle rotations?

$$\Omega(D) \sim \frac{1}{\tau(D)} \sim \frac{\delta_D u}{D} \sim D^{-2/3}$$
 (CGH)

• What is the effect of shape anisotropy on finite-sized particles?

Small anisotropic particles align with flow structures: For large particles is a randomization of rotations to be expected?

• What do we learn of turbulence when we look at a moving particle?

A new numerical study

in collaboration with

Chao Sun (孙超) & Linfeng Jiang (蒋林峰) Tsinghua University, Beijing, China



Navier-Stokes equations (NSE)



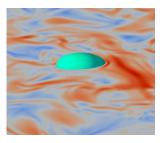
$$\partial_t \mathbf{u} + \mathbf{u} \cdot \nabla \mathbf{u} = -\rho^{-1} \nabla p + \nu \nabla^2 \mathbf{u} + \mathbf{f},$$

 $\nabla \cdot \mathbf{u} = 0,$

+ Newton-Euler equations (NEE)

$$m_{p} \frac{\mathrm{d}\mathbf{v}}{\mathrm{d}t} = \mathbf{F} \qquad \mathbf{F} = \oint_{S_{p}} \boldsymbol{\sigma} \cdot \mathbf{n} \, dS \qquad \boldsymbol{\sigma} = -p\mathbf{I} + \rho\nu(\nabla \mathbf{u} + \nabla \mathbf{u}^{T})$$
$$\frac{\mathrm{d}\boldsymbol{\mathcal{I}}\boldsymbol{\Omega}}{\mathrm{d}t} = \mathbf{T} \qquad \mathbf{T} = \oint_{S_{p}} (\mathbf{x} - \mathbf{r}) \times (\boldsymbol{\sigma} \cdot \mathbf{n}) \, dS$$

+ no-slip boundary conditions @fluid-particle interface



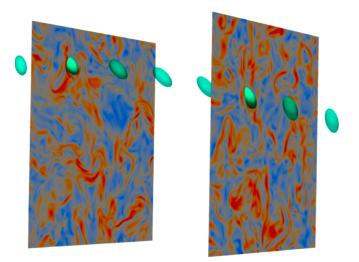
Code: <u>https://github.com/ecalzavarini/ch4-project</u>

Methods:

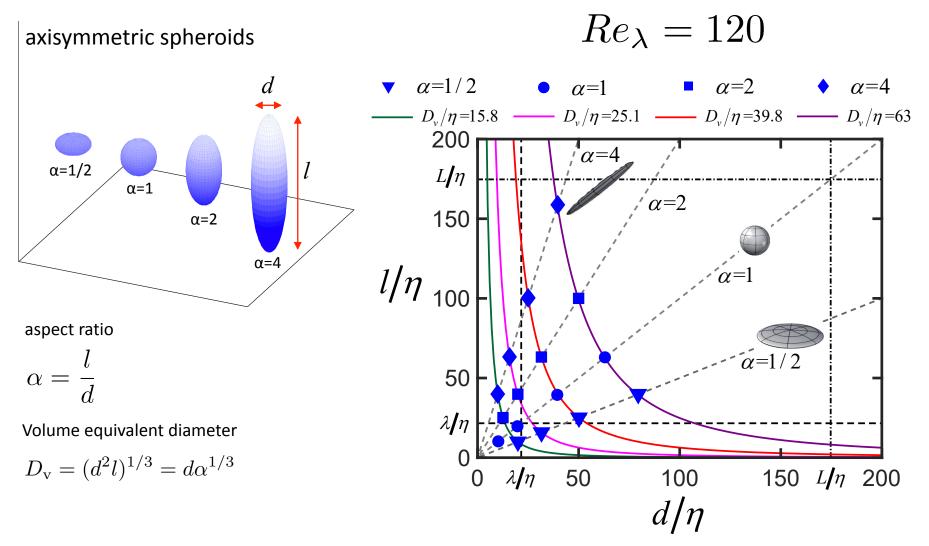
- Lattice Boltzmann (fluid)
- Quaternions (rigid body)
- Immersed boundaries (interface)

Same methods as

K. Suzuki, T. Inamuro, Comput. & Fluids 49, 173-187 (2011)



Parameter space



Auxiliary particle models

$$\mathbf{F} = \oint_{S_p} \boldsymbol{\sigma} \cdot \mathbf{n} \, dS$$
$$\mathbf{T} = \oint_{S_p} (\mathbf{x} - \mathbf{r}) \times (\boldsymbol{\sigma} \cdot \mathbf{n}) \, dS$$
Real Particles (RP)

Numerical Experiments

 Real Particles (RP)
 NSE + NEE + IBM

 Volume Averages (VA)
 NSE

 Virtual Particles (VP)
 NSE + NEE

 Real Particles with Fixed Locations (RPFL)
 NSE + EE + IBM

 Virtual Particles with Fixed Locations (VPFL)
 NSE + EE

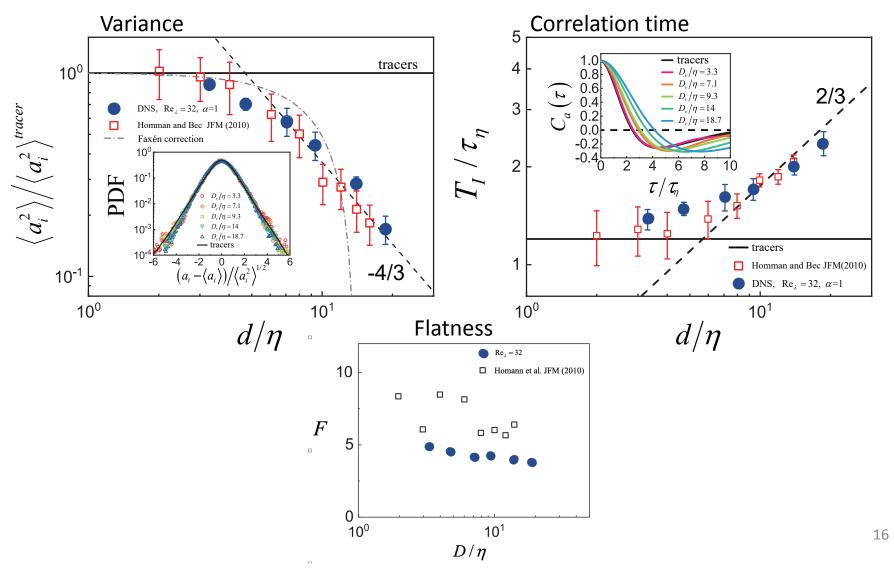
 Jeffery Fluid Tracers (JFT)
 NSE + tracer eq.+ Jeffery eq.

Model Equations

Spheres: acceleration at low Re

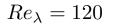
Comparison with Homann & Bec JFM 2010

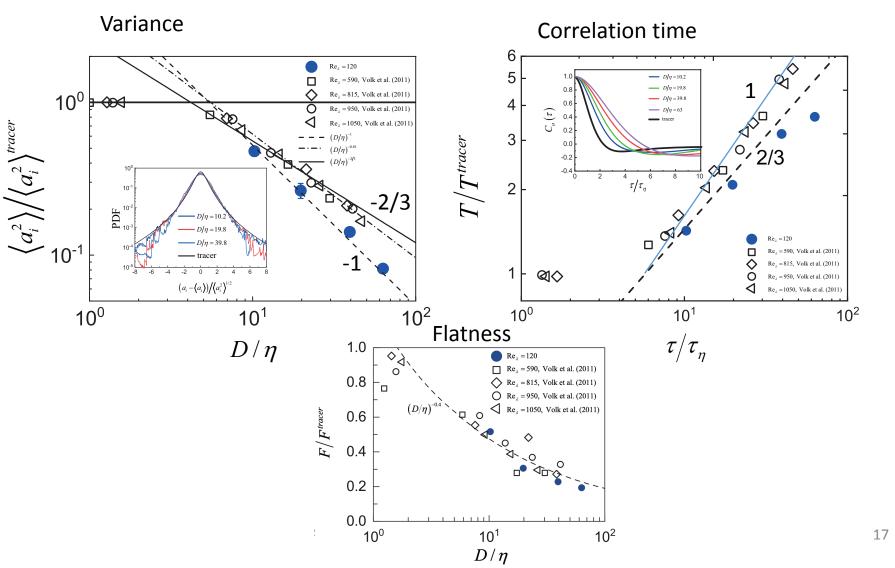




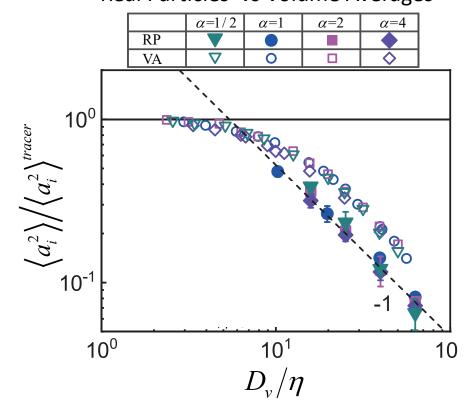
Spheres: acceleration at larger Re

Comparison with Volk et al. JFM 2011





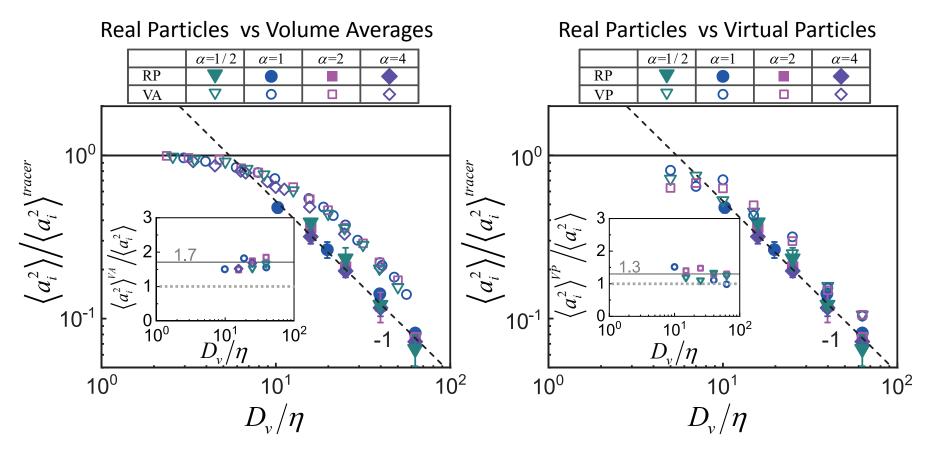
Spheroids: Acceleration variance



Real Particles vs Volume Averages

Same trend in D_v with an offset

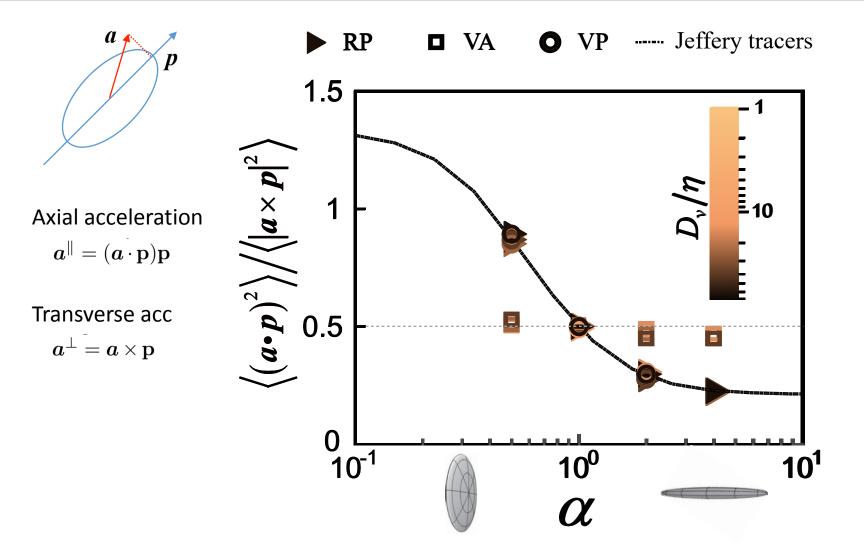
Spheroids: Acceleration variance



Better overlap for VP particles

$$D_{\rm v} = (d^2 l)^{1/3} = d\alpha^{1/3}$$

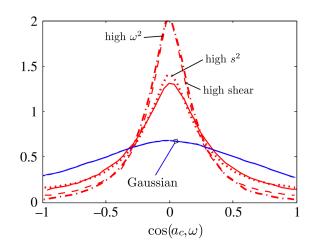
Spheroids: Acceleration alignment



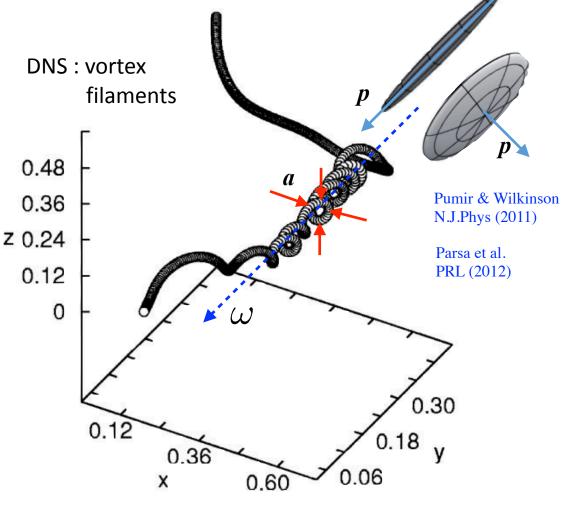
Spheroids: Acceleration alignment

A qualitative explanation

EXP: $\mathbf{a} \perp \boldsymbol{\omega}$



A.Liberzon et al. Physics (2012)



L.Biferale et al. Phys Fluids (2005)



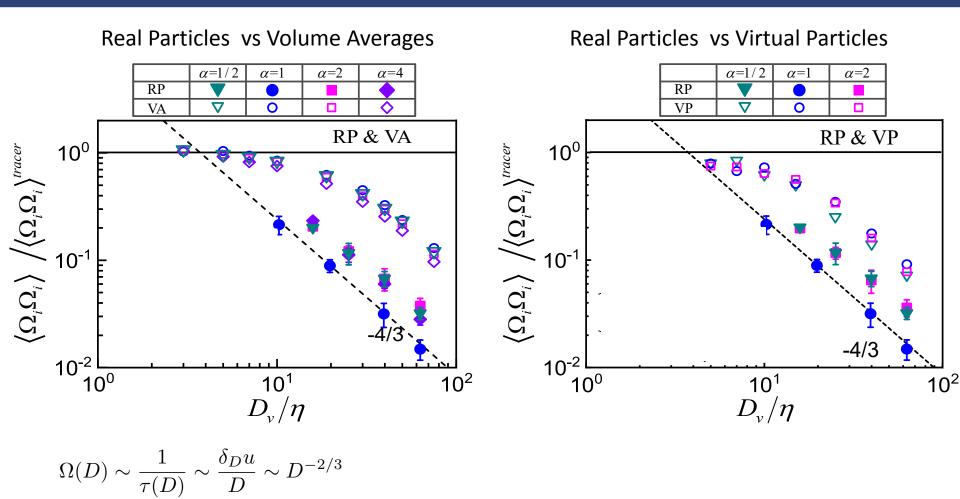
Acceleration

Size-dependence explained by fluid acceleration volume averages (VA or VP) Amplitude better accounted for dynamical (Lagrangian) volume averages (VP)

Importance of **flow structures**: preferential alignment even for inertial-scale particles!

Rotation

Spheroids: Angular velocity



But no overlap of RP with VA or VP

Spheroids: spinning and tumbling

Two predictions:

 $\langle \Omega_i^s \Omega_i^s \rangle = \frac{1}{12} \frac{1}{\tau_r^2}$

1) Randomly oriented model in HIT flow

Parsa, EC, Voth, Toschi, PRL 2012

Axial rotation rate (spin) $oldsymbol{\Omega}^s = (oldsymbol{\Omega} \cdot \mathbf{p}) \mathbf{p}$

 $\underline{\Omega}$

p

Transverse rotation (tumbling)

 $\dot{\mathbf{p}} = \boldsymbol{\varOmega} \times \mathbf{p}$

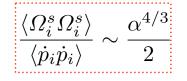
 $\langle \dot{p}_i \dot{p}_i \rangle = \left(\frac{1}{6} + \frac{1}{10}\Lambda^2\right) \frac{1}{\tau_\eta^2}$ $\Lambda = \frac{\alpha^2 - 1}{\alpha^2 + 1}$

2) CGH prediction

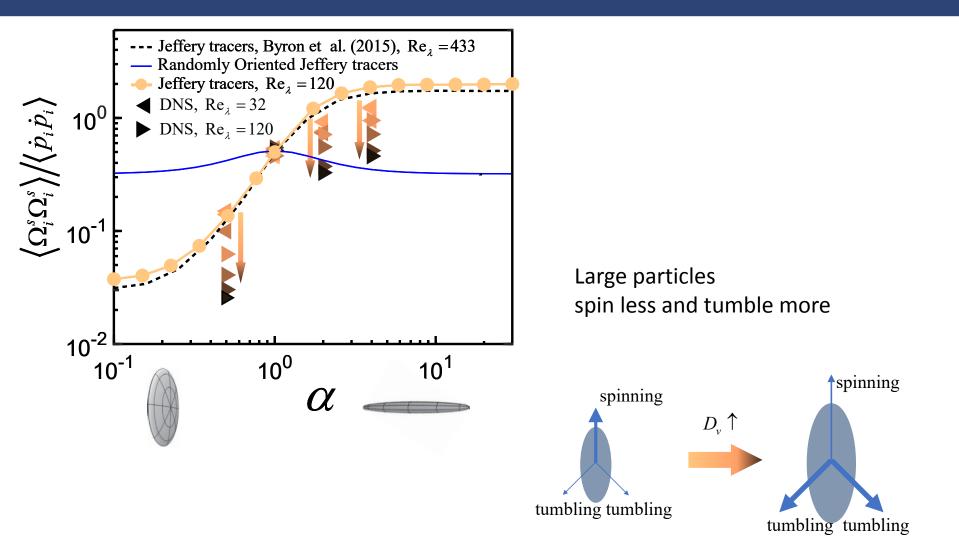
$$\begin{split} & \langle \Omega_i^s \Omega_i^s \rangle \sim d^{-4/3} \\ & \langle \dot{p}_i \dot{p}_i \rangle \sim l^{-4/3} \end{split}$$

$$\frac{\langle \Omega_i^s \Omega_i^s \rangle}{\langle \dot{p}_i \dot{p}_i \rangle} = \frac{\frac{1}{12} \langle \omega^2 \rangle}{\frac{1}{6} \langle \omega^2 \rangle + \frac{1}{5} \Lambda^2 \langle \mathcal{S} : \mathcal{S} \rangle} = \frac{5}{10 + 6\Lambda^2}$$

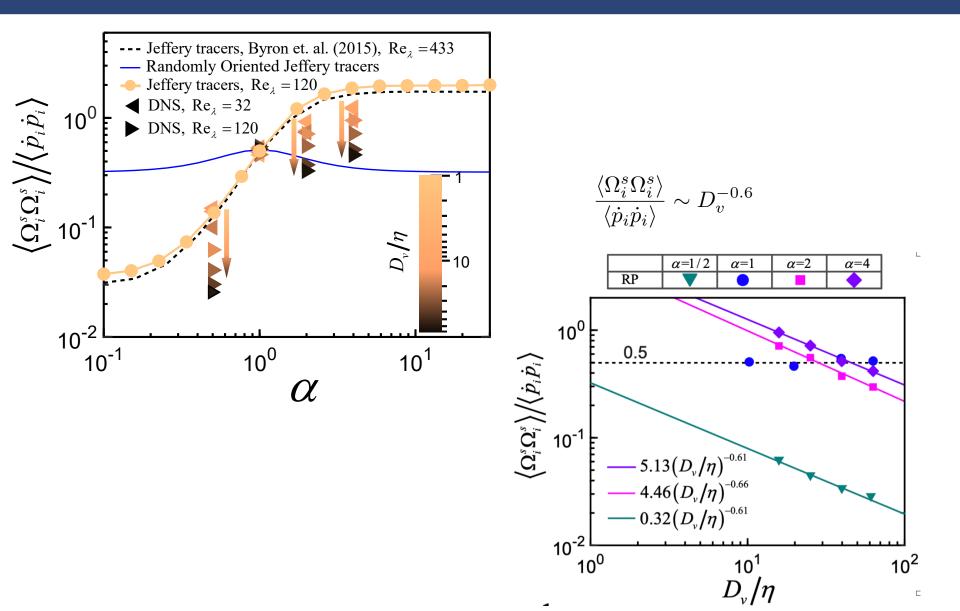
Oehmke et al., PRF 2021

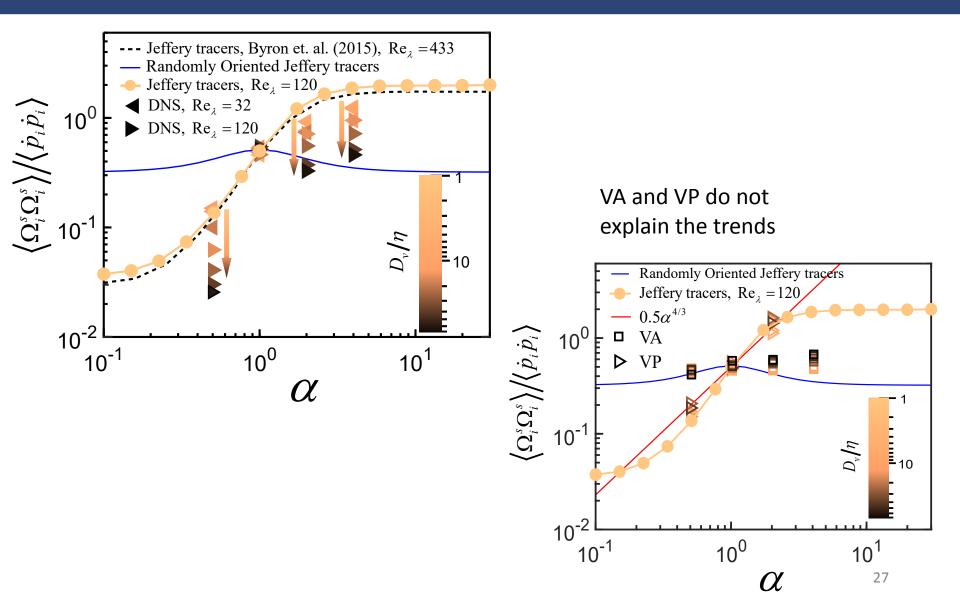


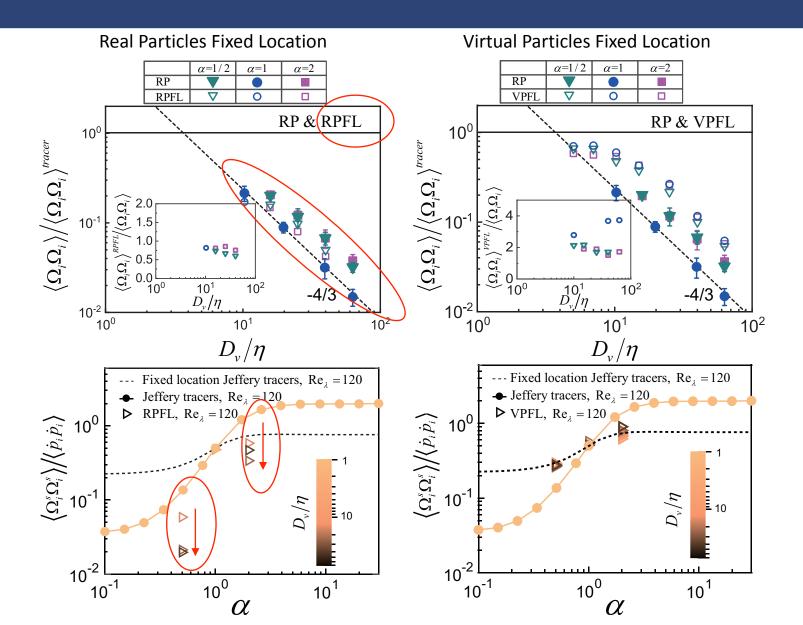
Always no size dependence!



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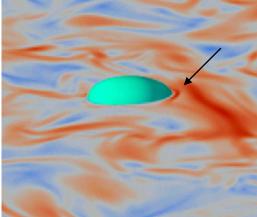




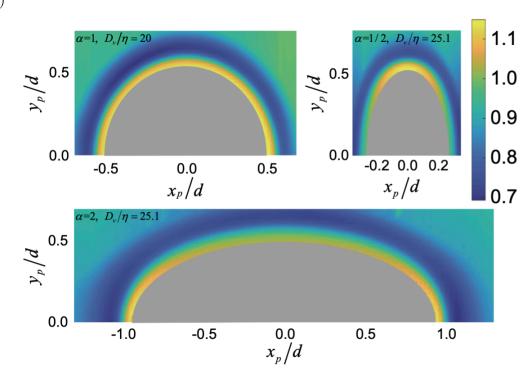
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A boundary layer effect?

$$Re_p = \frac{|\mathbf{u} - \mathbf{v}| D_v}{\nu} \simeq \frac{D_v}{\lambda} Re_\lambda \sim O(10^2$$



Dissipation field surrounding a fully resolved particle



 $\frac{1}{4}\overline{\boldsymbol{\omega}^2}/\overline{S:S}$

Overall increase in strain-rate as compared to vorticity -> this promote tumbling

Summary / conclusion

Acceleration

Size-dependence explained by fluid acceleration volume averages (VA or VP) Amplitude better accounted for by dynamical (Lagrangian) volume averages (VP)

Importance of **flow structures**: preferential alignment even for inertial-scale particles!

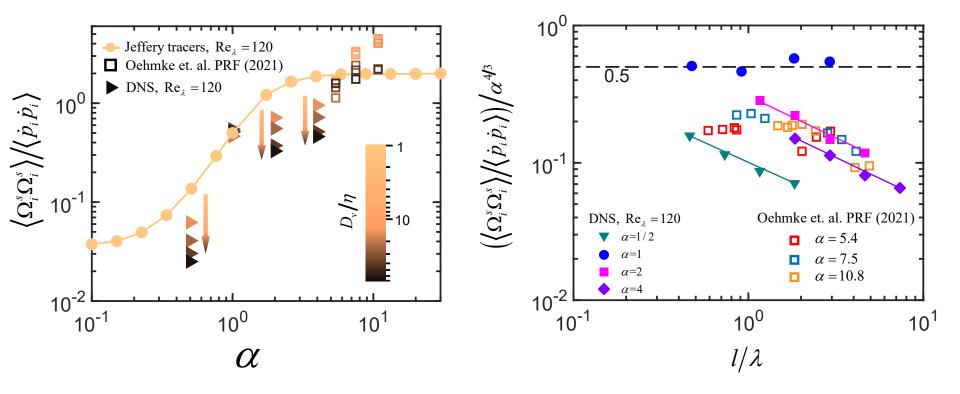
Rotation

Size-dependence explained in terms of volume averages but not its **amplitude** Importance of **two-way coupling** and **boundary layers**: decrease of spinning vs. tumbling!

Experimental verifications and high-resolution DNS are still (very much) needed!

Dynamics of finite-size spheroids in turbulent flow: the roles of flow structures and particle boundary layers, L. Jiang, C. Wang, S. Liu, C. Sun, E. Calzavarini, [http://arxiv.org/abs/2202.03937] **J. Fluid Mech** (2022)

Comparison with Oehmke, Bordoloi, Variano, Verhille PRF (2021)



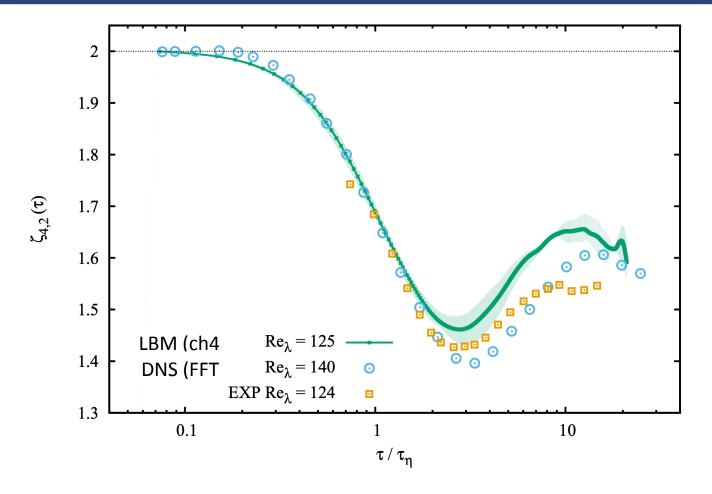
More experimental measures of tumbling & spinning needed.

Perspectives

- How to reach a more quantitative understanding? Modelling needed.
- Numerical results must be extended to higher Re numbers and to non-neutrally buoyant particles

Code validation: Lagrangian tracer statistics

Calz....PRF (2021)



$$\zeta_{4,2}(\tau) = \frac{d \log \langle (\delta_{\tau} \dot{x}_{s,i})^4 \rangle}{d \log \langle (\delta_{\tau} \dot{x}_{s,i})^2 \rangle},$$

Code validation: particle sedimentation

